



Diamond anvil cell behavior up to 4 Mbar HPSTAR 507-2018

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The diamond anvil cell (DAC) is considered one of the dominant devices to generate ultrahigh static pressure. The development of the DAC technique has enabled researchers to explore rich high-pressure science in the multimegabar pressure range. Here, we investigated the behavior of the DAC up to 400 GPa, which is the accepted pressure limit of a conventional DAC. By using a submicrometer synchrotron X-ray beam, double cupplings of the beveled diamond anvils were observed experimentally. Details of pressure loading, distribution, gasket-thickness variation, and diamond anvil deformation were studied to understand the generation of ultrahigh pressures, which may improve the conventional DAC techniques.

high pressure | diamond anvil cell | deformation

The diamond anvil cell (DAC) was invented almost 60 y ago (1). It squeezes samples in between two opposing diamonds to generate extremely high static pressure above 100 GPa (2). Many high-pressure DAC techniques have been developed to make it possible to generate pressures up to 400 GPa (3), such as the usage of a gasket, a beveled anvil design (4), different types of seat (5), etc. This makes the DAC a very important tool not only in geoscience, where it is used to study the earth's interior across the whole pressure range to the inner core ~360 GPa, but also for exploring the novel phenomena of matter with extremely high density in multidisciplinary scientific fields such as physics (6–8), chemistry (9–12), material science (13, 14), and even bioscience (15). Recently, an innovative use of a DAC with a pair of secondary nanodiamond anvils was reported to achieve above 6 Mbar (16, 17). However, its technical challenge and difficult sample preparation have limited its application as a routine way of generating high pressures.

In conventional DAC experiments, typically flat and beveled diamond anvils (Fig. 1, *Left*) are used to generate high pressures. For megabar pressure experiments (>100 GPa), beveled diamond anvils with small culets are usually recommended in practice. Many factors can cause the failure of diamond anvils, such as the DAC alignment, DAC stability, anvil quality, gasket selection, etc. Addressing all of these factors makes it challenging to achieve pressures above 300 GPa. Moreover, under a very high load, the beveled anvil will “cup” (bend outward), causing one of the major limitations on the achievable pressure for a beveled DAC (18). Thus, understanding the behavior of a DAC under very high pressures, including the process of pressure loading, the evolution of diamond anvil elastic deformation, gasket flow, and pressure distribution along the culet area, is indispensable for further understanding and developing the DAC technique.

Some previous experimental studies explored the behavior of the DAC under high pressures (18, 19). However, detailed local strain stress information (e.g., on the small culet area) could not be derived due to the limitation of the probe size, which typically has a 5 ~ 10- μ m spatial resolution. A few finite-element modeling studies also analyzed the anvil deformation (20, 21) and its stress state (21–23) under high pressures. For a beveled diamond anvil geometry, the finite-element modeling unveiled double cupping in the center flat and beveled areas at 300 GPa (21). However, this

double cupping has not been observed experimentally at high pressures up to 300 GPa. The experimental and theoretical discrepancy is likely due to either a low spatial resolution or not achieving high enough pressure in the experiments (18).

In this paper, by using submicrometer focused synchrotron X-ray beam, we performed in situ high-pressure synchrotron X-ray diffraction and absorption experiments to investigate the behavior of the DAC up to 400 GPa. The submicrometer X-ray offers many advantages in high-pressure research (24, 25), especially when dealing with small samples, because its high spatial resolution makes it possible to resolve details within a tiny culet area.

Mao-type symmetric DACs with culets of 20 μ m beveled to 250 or 200 μ m with a beveled angle of 8.5° on Boehler-type seats (Fig. 1, *Right*) were used to achieve 400-GPa pressure. To simplify the study, an initial thickness of 250- μ m tungsten (W) foil was used as both the sample and gasket. Pressure was applied by tightening the screws, thus bending the spring washers and then applying a force proportional to the angles of the screw rotation. The central pressure as a function of the screw rotation was recorded in Fig. 2*A* (solid dots). This pressure-loading curve is nonlinear and can be divided into three sections with different loading rates. We will discuss the loading curve in detail later. Fig. 2*B* presents a typical X-ray diffraction (XRD) pattern of W with a pressure of 170 GPa at the center of the culet. The insert shows the W-(110) peak shift during compression, which indicates a pressure increase from 170 to 398 GPa.

The pressure distribution as a function of the radial distance across the culet is shown in Fig. 3*A* and *C* at four selected center pressures of ~170, ~240, ~300, and ~400 GPa, respectively.

Significance

Diamond anvil cell (DAC) could be one of the most popular and versatile devices to generate static high pressure. By using beveled anvils, researchers can study materials up to 4 Mbar, which is considered as a limit pressure of conventional DAC; however, very few experiments have visited this pressure region. Here, we have systematically studied the behaviors of three DACs with three typical culet sizes, one of which reached 4 Mbar. In addition, by using submicrometer synchrotron X-ray probe, double cupping of beveled diamond anvils up to 4 Mbar is reported experimentally. Our results provide important guidance for future DAC studies to extreme high pressures.

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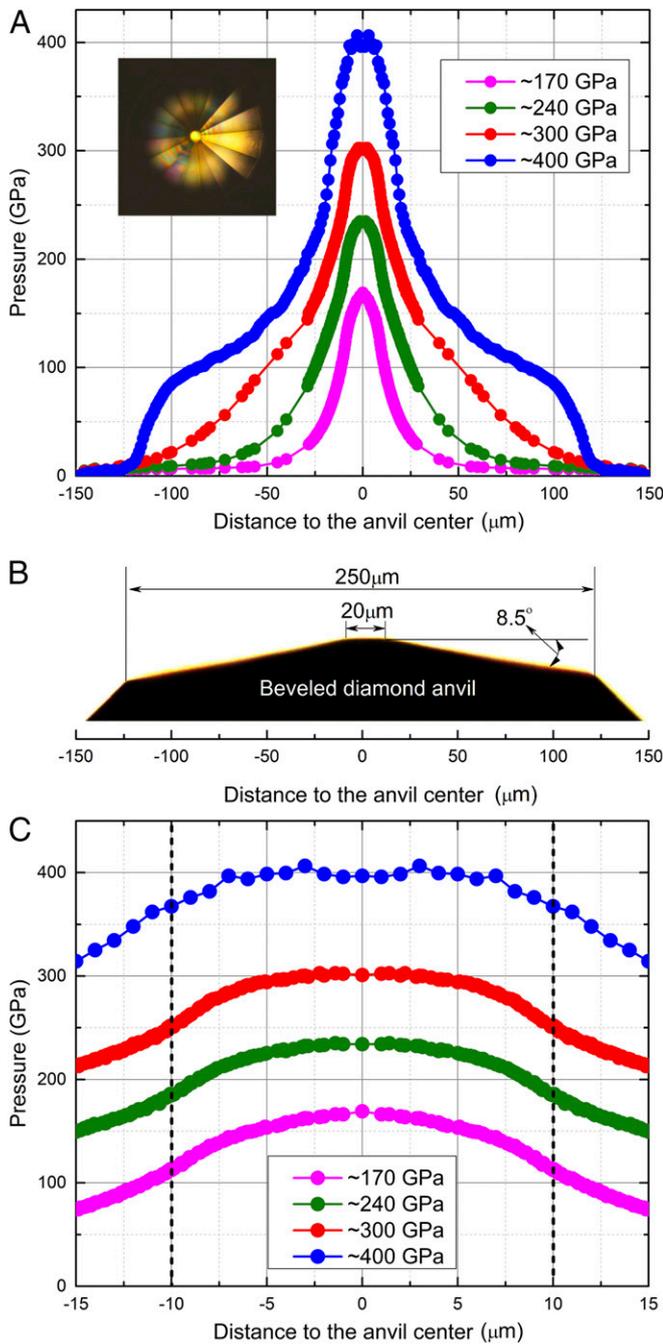


Fig. 3. Pressure distributions in a beveled DAC. Pressure distributions as a function of the radial distance from the center of the culet at four center pressures of ~170, ~240, ~300, and ~400 GPa, respectively. (A) Large region and (C) culet region. B is a side view of the diamond anvil with a 20- μm culet beveled to 250 μm at 8.5° as a reference. The photomicrograph in A shows the reflectance in a visual observation of the tiny culet and beveled area at ~400 GPa.

connecting the peak pressures at the center and edge pressures at $\pm 150 \mu\text{m}$ are presented as references for each center pressure. At center pressures below ~40 GPa, the pressure distribution is more or less linear from the center to the edge. However, when pressure is increased to above ~50 GPa, deviations from the linear pressure distribution appear and become increasingly severe up to ~90 GPa (the highest pressure in this study). A linear pressure distribution has been observed in many previous studies and is used to estimate the maximum shear stress (26) using the

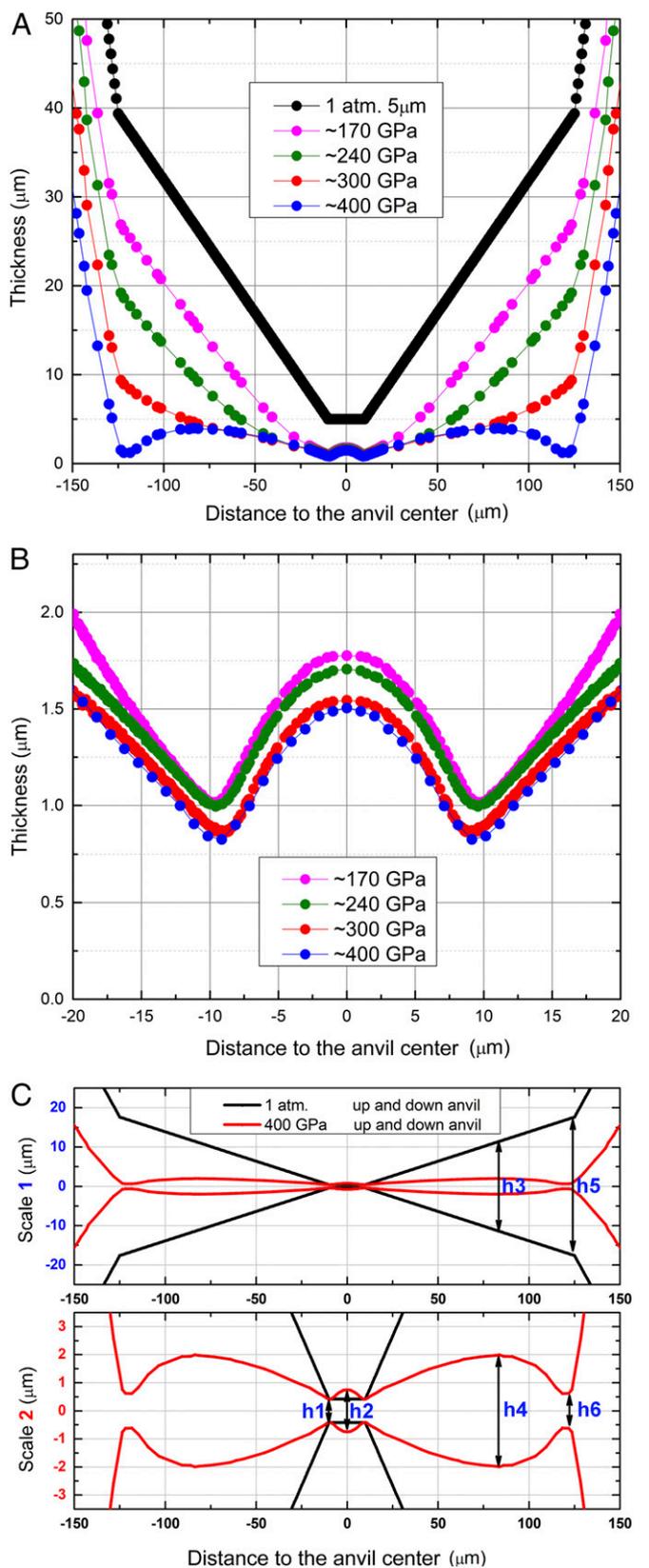


Fig. 4. W gasket-thickness distribution at various pressures and anvil shapes compared between 1 atm and ~400 GPa. Thickness distribution as a function of the radial distance from the center of the culet at four center pressures of ~170, ~240, ~300, and ~400 GPa, respectively. (A) Large region and (B) culet region. (C) Comparison of the anvil shapes between ambient pressure (black) and at ~400 GPa (red) in different scales, large (C, Upper) and small (C, Lower).

Table 1. Forces and average pressures at four different center pressures

Center pressure, GPa	Force on 250- μm culet, N	Force on 20- μm culet, N	Average pressure, GPa	Ratio of center/average pressure
170	540	45	11	15
240	1,126	67	23	10
300	2,612	89	53	6
400	4,901	122	100	4

equation $dP/dr \sim 2\tau/h$, where P is the pressure, r is the radius from the culet center, τ is the maximum shear stress, and h is the thickness of the sample. However, estimations using this method are inapplicable when pressure deviates significantly from a linear distribution. Fig. 5B shows the gasket-thickness variations as the pressure increases, from which anvil deformation can be inferred. First, the anvil cupping begins at very low pressure, although the bending of the culet surface is very slight. Second, the gasket thickness at the center does not change much after ~ 50 GPa, where the cupping becomes increasingly severe.

Fig. 5C shows the center pressure as a function of screw rotation, which is proportional to the loading force. The zero-angle positions from the experiments in both Figs. 2A and 5C were set where the screws were finger tight, then the angle of the screw rotation and the center pressure were recorded when DACs were loaded. The loading curve can be divided into three stages with two intersectional points at about 10 and 50 GPa. In stage I, pressure increased slowly with plastic deformation of the W gasket; in this stage, the gasket flows significantly to release the force accumulated on the culet with a slow pressure increase. In stage II, when a linear pressure distribution has been established from ~ 10 to ~ 50 GPa, the gasket becomes “locked” and is gradually flattened with a rapid center pressure increase. Finally, in stage III above ~ 50 GPa, pressures are redistributed, and total “cupping” presents. In this stage, a portion of the force has to be used to help pressure redistribute and not only work to increase the peak pressure in the center. Thus, the loading rate is reduced more slowly than in stage II.

The forces on the culet area can also be calculated like those in Fig. 2A with different pressure distributions at each center pressure. Fig. 5C shows the center pressure as a function of calculated forces on the 300- μm culet (open circles). The curve from the calculation matches stages II and III very well, which confirms that the loading-rate changes are the result of pressure redistribution between stages II and III. The calculated forces on the culet area in Figs. 2A and 5C are part of the forces totally applied to the DACs. In real experiment, there are many other forces such as the one applying on the pavilion part due to the gasket extruding out of the culet, the one to bend the DAC itself, the one against frictions during load, etc. However, the force on the culet area is a major contribution of the applied force to the DAC, especially when the DAC is under relatively high pressures.

The pressure-load curve with 300- μm flat culets has the same “s” shape as the one using 20- μm beveled to 250- μm culets. The s-shaped curve in the 400-GPa experiment (Fig. 2A) has the same features: a slower loading rate in stages I and III than in stage II, which is caused by gasket plastic deformation and pressure redistribution, respectively. A similar experiment on 100- μm -beveled to 300- μm -beveled culets was also performed up to 175 GPa, and an s-shaped curve was also observed (Fig. S1). This s-shape loading behavior is universal in all conventional DAC experiments.

In summary, we studied the behaviors of conventional DACs up to 400 GPa using a submicrometer X-ray beam. The relationship between the s-shaped load curve, pressure distribution, gasket thickness variation, and anvil deformation on both the beveled

and flat anvils were investigated in detail. This study shows that an s-shaped loading behavior and unavoidable cupping in both beveled and flat anvils are universal in conventional DAC experiments.

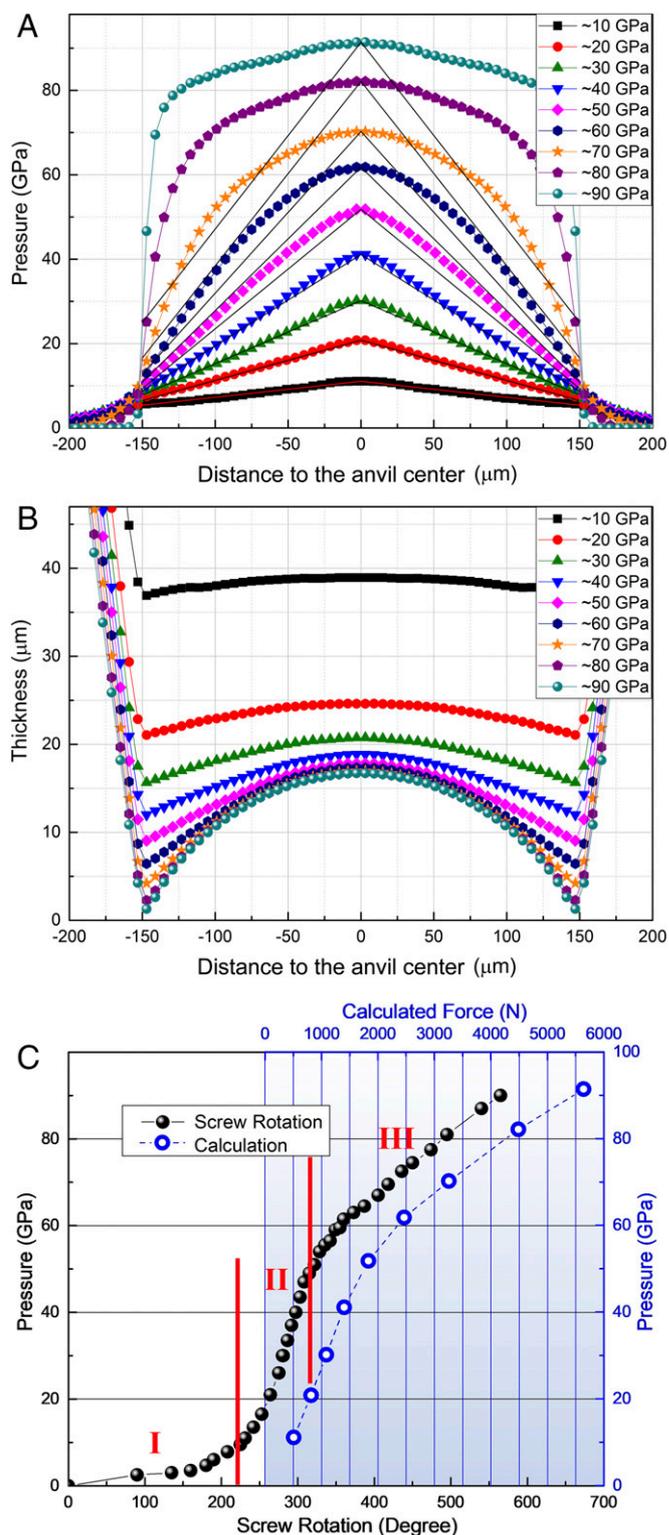


Fig. 5. The behavior of a flat 300- μm -anvil DAC up to 90 GPa. Pressure distributions (A) and gasket-thickness variations (B) as a function of the radial distance from the center of the culet at different center pressures. (C) The pressure load curves, solid dots, and open circles are from the experiment and calculations, respectively.

In addition, double cupping up to 400 GPa was observed experimentally both in a small culet area and the whole beveled-anvil area. The small culet area first cups and then remains roughly the same up to the highest load, while the whole beveled-anvil area experiences continuous cupping. To exceed the pressure limit of a conventional DAC, using a superhard nanocrystalline diamond (27, 28) up to the anvil material could be a solution (29).

Methods

High pressures were generated using Mao-type symmetric DACs with Boehler-type seats. Beveled anvils with culet size of 20 μm beveled to 250 or 200 μm ; 100 μm beveled to 300 μm with all beveled angle of 8.5°; and flat anvils with culet size of 300 μm were used to achieve \sim 400-, \sim 175-, and \sim 90-GPa pressures, respectively. All of the experiments were performed at room temperature.

In situ high-pressure XRD was carried out to monitor the pressure changes and distribution on both the culet and beveled areas at beamlines 34-ID-E, 2-ID-D, and 16-ID-B of the Advanced Photon Source, Argonne National Laboratory, and the BL15U1 beamline, Shanghai Synchrotron Radiation Facility in China. Diffraction patterns were recorded using Mar-CCD or Pilatus area detectors. The X-ray beam sizes at 34-ID-E, 2-ID-D, 16-ID-B, and BL15U1 were \sim 500 nm, \sim 250 nm, $1 \times 2 \mu\text{m}$, and $2 \times 2 \mu\text{m}$, respectively.

To measure the gasket thickness, in situ high-pressure X-ray absorption profiles were measured at beamlines 34-ID-E and 16-ID-B. According to absorption law, the equation $I/I_0 = \exp(-\mu\rho d)$ was used to calculate the thickness of the W gasket, where I , I_0 , μ , ρ , and d represent the X-ray intensity after the sample, incident X-ray intensity before the sample, mass attenuation coefficient, density, and thickness of the sample (W), respectively. The density of W was derived from its equation of state (30) using the XRD data. The mass attenuation coefficient of W was obtained from National Institute of Standards and Technology (31). A normalized transmission intensity through the empty DAC without the W gasket was used as I_0 in this study; and the normalized transmission intensity through the W gasket and DAC was used as I . For the 400-GPa experiment, two X-ray energies of 14 and 21 keV were selected to optimize the contrast of the absorption measurement considering the comparison between the thickness and the absorption length of the W gasket. Typically, at low pressures when

the gasket was thicker, a 21-keV X-ray was used; for a culet area where the gasket was thin, a 14-keV X-ray was always used. A 30.5-keV X-ray was used for the experiments to \sim 175 and \sim 90 GPa.

For both the XRD and X-ray absorption measurement in the 400-GPa experiment, typically a 2D X-ray absorption mapping (\sim 1-h scanning time) was first measured after the pressure was increased to the target pressure, and then a 2D XRD mapping (\sim 5-h scanning time) was performed. The 2D mappings were performed at selected pressures to cover the whole culet area with two different step sizes: (i) Large-area scan: step size 20 μm , coverage $320 \times 320 \mu\text{m}^2$. (ii) Small-area scan: step size 1 μm , coverage $40 \times 40 \mu\text{m}^2$. Both mappings were centered on the culet center. Circular symmetry was assumed for the 2D maps, and the data were converted into radial 1D data for averaging (Fig. S2A). One-dimensional XRD scan on culet of 20 μm beveled to 200 μm was performed to get the pressure distribution at the center pressure of \sim 400 GPa. Pressure distribution on 20 μm beveled to 200 μm was then converted to that of 20 μm beveled to 250- μm culet. The 20- μm area was kept the same but the beveled part from 20 to 200 μm was scaled up to 250 μm . For the experiments to \sim 175 and \sim 90 GPa, 1D line-scan data were collected and reflection symmetry for averaging was used for both the XRD and X-ray absorption measurements.

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