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# Charge and Spin-State Characterization of Cobalt Bis(*o*-dioxolene) Valence Tautomers Using Co $K\beta$ X-ray Emission and L-Edge X-ray Absorption Spectroscopies

H. Winnie Liang,<sup>\*,†,‡,§</sup> Thomas Kroll,<sup>§</sup> Dennis Nordlund,<sup>§</sup> Tsu-Chien Weng,<sup>||</sup> Dimosthenis Sokaras,<sup>§</sup> Cortlandt G. Pierpont,<sup>#</sup> and Kelly J. Gaffney<sup>\*,†,‡,§</sup>

<sup>†</sup>Department of Chemistry, Stanford University, Stanford, California 94305, United States

<sup>‡</sup>PULSE Institute, SLAC National Accelerator Laboratory, Menlo Park, California 94025, United States

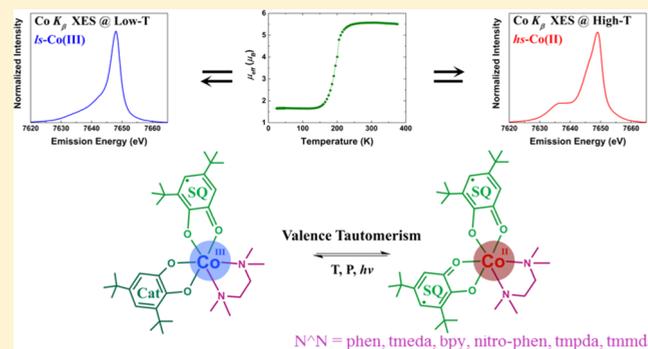
<sup>§</sup>Stanford Synchrotron Radiation Lightsource (SSRL), SLAC National Accelerator Laboratory, Menlo Park, California 94025, United States

<sup>||</sup>Center for High Pressure Science & Technology Advanced Research, Pudong, Shanghai 201203, P. R. China

<sup>#</sup>Department of Chemistry and Biochemistry, University of Colorado, Boulder, Colorado 80309, United States

## Supporting Information

**ABSTRACT:** The valence tautomeric states of Co(phen)-(3,5-DBQ)<sub>2</sub> and Co(tmeda)(3,5-DBQ)<sub>2</sub>, where 3,5-DBQ is either the semiquinone (SQ<sup>-</sup>) or catecholate (Cat<sup>2-</sup>) form of 3,5-di-*tert*-butyl-1,2-benzoquinone, have been examined by a series of cobalt-specific X-ray spectroscopies. In this work, we have utilized the sensitivity of 1s3p X-ray emission spectroscopy ( $K\beta$  XES) to the oxidation and spin states of 3d transition-metal ions to determine the cobalt-specific electronic structure of valence tautomers. A comparison of their  $K\beta$  XES spectra with the spectra of cobalt coordination complexes with known oxidation and spin states demonstrates that the low-temperature valence tautomer can be described as a low-spin Co<sup>III</sup> configuration and the high-temperature valence tautomer as a high-spin Co<sup>II</sup> configuration. This conclusion is further supported by Co L-edge X-ray absorption spectroscopy (L-edge XAS) of the high-temperature valence tautomers and ligand-field atomic-multiplet calculations of the  $K\beta$  XES and L-edge XAS spectra. The nature and strength of the magnetic exchange interaction between the cobalt center and SQ<sup>-</sup> in cobalt valence tautomers is discussed in view of the effective spin at the Co site from  $K\beta$  XES and the molecular spin moment from magnetic susceptibility measurements.



## INTRODUCTION

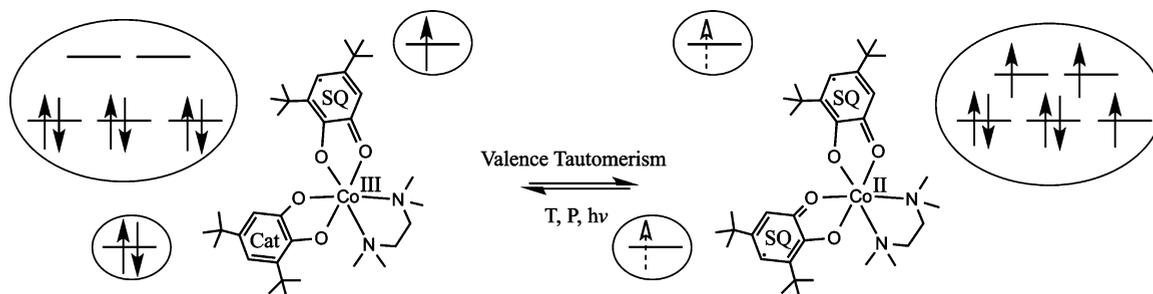
The ability to generate long-lived metastable electronic states with external stimuli makes the phenomena of spin crossover (SCO) and valence tautomerism (VT) in 3d transition-metal complexes a subject of great interest, with potential applications in spintronic devices, data storage, and molecular switches.<sup>1–6</sup> The thermodynamic ease with which coordination complexes with redox-active ligands can convert between distinct charge and spin states has also been associated with catalytic activity in organic transformations.<sup>7–9</sup> VT corresponds to an intramolecular electron transfer between a redox-active metal center and a redox-active ligand coupled to a change in spin multiplicity, or a spin crossover, at the metal center.<sup>10–16</sup> Furthermore, as is true for SCO transitions, VT transitions are typically accompanied by distinct and reversible changes in the structural, optical, and magnetic properties of the metal complexes.<sup>15,17–19</sup> They have been observed in transition-metal complexes of Mn, Fe, Co, Ni, Cu, Rh, Ru, and Ir with a

variety of redox-active ligand types, including *o*-dioxolenes, porphyrins, phenoxy ligands, ferrocenes, and tetrathiafulvalene.<sup>12,20–27</sup> In general, metal complexes displaying VT transitions meet the following two electronic requirements: (1) the degree of covalency in the interaction between a metal ion and a redox-active ligand must be small, and (2) their frontier orbitals must be of similar energy.<sup>2,12</sup>

Cobalt bis(*o*-dioxolene) coordination compounds represent an archetypical series of valence tautomeric complexes, where the term “*o*-dioxolene” denotes the redox series of neutral *o*-benzoquinone (*o*-BQ), monoanionic semiquinone (SQ<sup>-</sup>), and catecholate dianion (Cat<sup>2-</sup>), without specifying the substituents on the benzene ring. In [Co<sup>III</sup>(3,5-DBSQ)(3,5-DBCat)(N<sup>N</sup>)], where 3,5-DBSQ/3,5-DBCat is a 3,5-di-*tert*-butyl-substituted semiquinone/catecholate and N<sup>N</sup> is a bidentate N-donor

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**Scheme 1. Illustration of VT in  $\text{Co}(\text{tmeda})(3,5\text{-DBQ})_2$  with the Co 3d and Ligand  $\pi^*$  Orbitals Shown, Where DBQ Can Either Be a Singly-Reduced Semiquinone ( $\text{SQ}^-$ ) or a Doubly-Reduced Catecholate Ligand ( $\text{Cat}^{2-}$ ) Depending on the Valence Tautomeric State of the Complex at the Specified Temperature, with the  $\text{Co}^{\text{III}}(\text{tmeda})(3,5\text{-DBSQ})(3,5\text{-DBCat})$  Form Favored at Low Temperature and the  $\text{Co}^{\text{II}}(\text{tmeda})(3,5\text{-DBSQ})_2$  Form Favored at High Temperature<sup>a</sup>**



<sup>a</sup>The redox series of *o*-BQ can be described as  $\text{Cat}^{2-} \rightarrow \text{SQ}^- + e^- \rightarrow o\text{-BQ} + 2e^-$ . Note that the coordination geometry at the Co center is assumed to be pseudooctahedral, with the metal d level split into the  $t_{2g}$  and  $e_g$  sets, whereas the dashed electron arrows on  $\text{SQ}^-$  are a reminder of the weak coupling interactions between unpaired electrons on the ligands and those at the metal center.

coligand such as 1,10-phenanthroline (phen) and  $N,N,N',N'$ -tetramethylethylenediamine (tmeda), intervalence electron transfer nominally involves the transfer of an electron from the  $\text{Cat}^{2-}$   $\pi^*$  orbital to the low-spin (ls)- $\text{Co}^{\text{III}}$  center to form a  $\text{SQ}^-$  and a high-spin (hs)  $\text{Co}^{\text{II}}$  (Scheme 1). It has been shown that valence tautomeric interconversion can be driven by changes in the temperature ( $T$ ) and pressure ( $P$ ), as well as by irradiation with lasers or soft X-rays.<sup>28–37</sup> In this work, all cobalt valence tautomers will be referred to as either  $\text{Co}(\text{N}^{\wedge}\text{N})$  or  $\text{Co}(\text{N}^{\wedge}\text{N})(3,5\text{-DBQ})_2$ , where  $\text{N}^{\wedge}\text{N}$  specifies the coligand identity; DBQ can refer to either semiquinone or catecholate, depending on the valence tautomeric state of the complex at the specified temperature.

Our study addresses the number of electronic configurations needed to accurately describe the electronic structure of the low- and high- $T$  valence tautomeric states, an open question despite extensive studies of cobalt bis(*o*-dioxolene) valence tautomers. In 2002, LaBute et al.<sup>38,39</sup> used a variational configuration interaction approach and Co K-edge X-ray absorption spectroscopy (K-edge XAS) to investigate the  $\text{Co}(\text{phen})$  system. They concluded that multiple configurations must be included in the calculations to accurately reproduce the energy difference between the low- and high- $T$  states and the 1s3d preedge absorption features in K-edge XAS, with the high- $T$  valence tautomeric phase having both hs- $\text{Co}^{\text{II}}$  and hs- $\text{Co}^{\text{III}}$  character and the low- $T$  phase having both ls- $\text{Co}^{\text{II}}$  and ls- $\text{Co}^{\text{III}}$  character. An additional theoretical study by Sato et al.<sup>40</sup> also highlights the potential significance and energetic accessibility of a ls- $\text{Co}^{\text{II}}$  configuration in VT. On the basis of their study of the minimum-energy crossing points and the potential energy paths connecting distinct spin states of cobalt, they proposed a mechanism in which the ls- $\text{Co}^{\text{II}}$  species can serve as an intermediate state in the thermally driven interconversion of ls  $\text{Co}^{\text{III}}$  to hs  $\text{Co}^{\text{II}}$ .

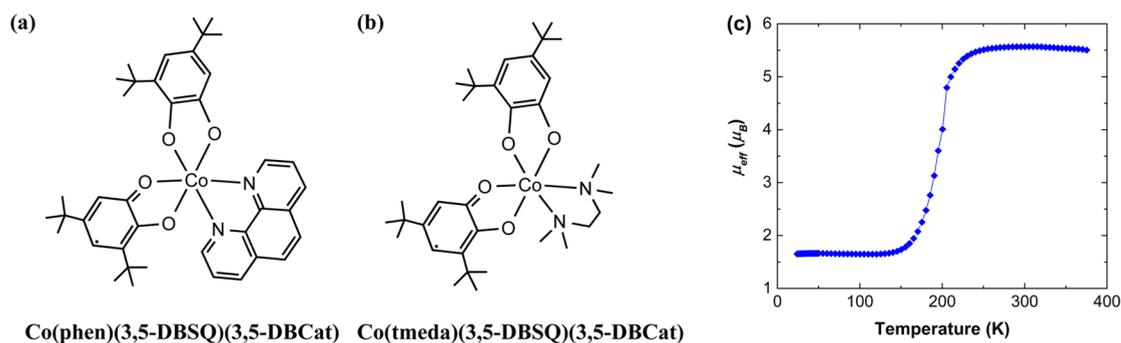
In this work, we present a Co 1s3p X-ray emission spectroscopy ( $K\beta$  XES) and Co L-edge X-ray absorption spectroscopy (L-edge XAS) investigation of the electronic structures of the  $\text{Co}(\text{phen})$  and  $\text{Co}(\text{tmeda})$  valence tautomers. Both of these spectroscopies have been well demonstrated to characterize the charge, spin, and metal–ligand covalency of 3d coordination complexes from the metal’s perspective, which will be used to investigate the importance of multiconfigurational interactions and the strength of the  $\text{Co}\text{--}\text{SQ}^-$  magnetic exchange interaction in valence tautomers. The atomic

specificity of X-ray spectroscopies differentiates them from conventional techniques such as electron paramagnetic resonance (EPR), UV–visible, IR, NMR, and magnetic susceptibility. While these non-X-ray-based methods allow the interconversion process to be monitored in a temperature-dependent manner, they lack the capability to discern changes in the spin and charge configurations at the Co center. The strong exchange interaction between the unpaired 3p and 3d electrons in the final state created by X-ray emission controls the relative intensity and energetic positions of the  $K\beta_{1,3}$  and  $K\beta'$  spectral features in the XES spectrum and makes the measurement sensitive to the effective metal spin and charge states, which can be further modulated by metal–ligand covalency.<sup>41–44</sup> We have also used Co L-edge XAS measurements and ligand-field atomic-multiplet calculations to investigate the cobalt-specific electronic structure at high temperature. These measurements and calculations complement  $K\beta$  XES and enable a more comprehensive determination of the importance of a hs- $\text{Co}^{\text{III}}$  configuration in the high- $T$  valence tautomers. Both XES and XAS measurements, as well as atomic-multiplet calculations, demonstrate that both valence tautomeric states can be robustly described with single-electron configurations for cobalt and provide no evidence for the hs- $\text{Co}^{\text{III}}$  configuration proposed by previous researchers.<sup>39</sup> During our measurements, we also observed the occurrence of soft X-ray-induced VT in our low- $T$  L-edge XAS data that is consistent with the work of Poneti et al.<sup>32</sup> and the possibility of hard X-ray-induced VT with aromatic nitrogen coligands. Given prior reports of interconversions induced by hard X-rays,<sup>45–50</sup> this, in turn, has limited our XES study of the electronic structure of low- $T$  valence tautomers to the  $\text{Co}(\text{tmeda})$  complex.

By combining the cobalt-specific charge and spin states derived from Co  $K\beta$  XES with the molecular spin moment extracted from magnetic susceptibility measurements, we proceed to address the strength of the magnetic exchange interaction between unpaired spins on the  $\text{SQ}^-$  ligands and the Co center in high- $T$  valence tautomers. These measurements show a spin moment of  $S = 3/2$  for both the Co ion and the entire molecule, an observation consistent with weak magnetic exchange interactions between  $\text{SQ}^-$  and the hs  $d^7$  metal center.

## EXPERIMENTAL SECTION

**Sample Preparation.** ls- and hs- $\text{Co}^{\text{III}}$  and - $\text{Co}^{\text{II}}$  model compounds were purchased from Sigma-Aldrich and used as they were received:



**Figure 1.** Cobalt-containing valence tautomeric complexes investigated in this work: (a)  $\text{Co}(\text{phen})(3,5\text{-DBQ})_2\cdot\text{C}_6\text{H}_5\text{CH}_3$  and (b)  $\text{Co}(\text{tmeda})(3,5\text{-DBQ})_2\cdot 0.5\text{C}_6\text{H}_6$ . (c) Temperature-dependent effective magnetic moment of  $\text{Co}(\text{tmeda})(3,5\text{-DBQ})_2\cdot 0.5\text{C}_6\text{H}_6$ . Molecular structures are expressed in their low- $T$  valence tautomeric state. Note that Q can be either semiquinone or catecholate, depending on the valence tautomeric state of the complex at the specified temperature.

tris(ethylenediamine)cobalt(III) chloride dihydrate, cobalt(III) acetylacetonate, and cobalt(II) acetylacetonate, namely,  $[\text{Co}(\text{en})_3]\text{Cl}_3\cdot 2\text{H}_2\text{O}$ ,  $\text{Co}(\text{acac})_3$ , and  $\text{Co}(\text{acac})_2$ , respectively. Compounds (a)  $\text{Co}(\text{phen})(3,5\text{-DBQ})_2\cdot\text{C}_6\text{H}_5\text{CH}_3$  and (b)  $\text{Co}(\text{tmeda})(3,5\text{-DBQ})_2\cdot 0.5\text{C}_6\text{H}_6$  (see Figure 1), along with  $\text{Co}_4(3,5\text{-DBSQ})_8$  and  $\text{Co}(\text{phen})_3\cdot\text{Cl}_2\cdot 4\text{H}_2\text{O}$ , were synthesized according to literature procedures<sup>51–54</sup> with reagents and solvents purchased directly from Sigma-Aldrich. The final products were characterized in toluene by UV–visible and IR absorption spectroscopies and confirmed with elemental analysis. Samples for X-ray emission experiments were ground into a fine powder, packed tightly into a 1.0-mm-thick aluminum cell, and sealed with 63.5  $\mu\text{m}$  kapton tape. For air-sensitive samples or samples going into a vacuum, as in soft X-ray measurements, the preparation step was done inside a glovebox under an inert  $\text{N}_2$  atmosphere.

**XES Data Collection and Processing.**  $K\beta$  X-ray emission spectra were collected at SSRL beamline 6-2 with ring conditions of 3.0 GeV and 500 mA. The incident beam energy was set to 8.0 keV using a  $\text{LN}_2$ -cooled Si(111) monochromator that was calibrated with a cobalt foil, providing  $\sim 2 \times 10^{13}$  photons  $\text{s}^{-1}$  in a  $0.15 \times 0.4$  mm spot size and an energy resolution ( $\Delta E/E$ ) of  $1 \times 10^{-4}$ . Emitted X-rays were energy-selected using four spherically bent Ge(444) crystals arranged in a Rowland geometry and were detected with a silicon drift detector. The incident flux was measured using a helium-filled ion chamber just upstream of the sample. A helium-filled flight path was used between the sample, crystal analyzers, and detector to minimize signal attenuation. X-ray-induced damage was assessed by collecting successive X-ray scans on a single spot; multiple spots were used when needed for each sample. Temperature control of the samples was achieved using a combination of a helium cryostat and a resistive heater.

In all of the samples, two scans were taken per spot, for a total of six scans per sample. These scans were then each normalized by its incident flux and averaged together to yield the final spectrum in PyMca. Each average spectrum was further scaled by its total intensity under the curve to allow for comparison.

**L-Edge XAS Data Collection and Processing.** Co L-edge X-ray absorption spectroscopy was performed at SSRL beamline 10-1. The beamline is equipped with a 6 m spherical grating monochromator operated with a 1000 lines  $\text{mm}^{-1}$  grating (Co 2p XAS) and 20  $\mu\text{m}$  entrance and exit slits, providing  $10^{11}$  photons  $\text{s}^{-1}$  at a resolution  $\Delta E/E$  of  $2 \times 10^{-4}$  in a beam spot of 1.0  $\text{mm}^2$ .

All data were acquired in the temperature range of 90–298 K, where the samples were cooled by  $\text{LN}_2$  and under an ultrahigh vacuum ( $\sim 10^{-9}$  Torr). Both the total electron yield (TEY, measured via the sample drain current, SC) and Auger electron yield (AEY, measured with a cylindrical mirror analyzer coupled to a channeltron) were recorded simultaneously. All spectra were normalized by the incoming flux, measured from a thin grid with freshly evaporated gold, positioned upstream of the sample chamber. A reference sample with several transition-metal oxides, intercepting a few percent of the beam upstream and recorded simultaneously, was used to calibrate the

energy of the cobalt scans with a relative energy precision of 25 meV. All powder samples were thoroughly mixed with 50% graphite by volume and affixed to an aluminum sample holder using conductive carbon tape. The AEY was used for the XAS spectra in this study. Background correction was performed using the statistics-sensitive nonlinear iterative peak-clipping algorithm in PyMca. Each spectrum was an average of 4–5 scans, with each scan taken at a different spot on the sample tape.

**Ligand-Field Atomic-Multiplet Calculations.** Atomic-multiplet calculations were performed using the CTM4XAS 5.5 program.<sup>55,56</sup> The effects of covalent metal–ligand mixing were included as a screening of the 3d–3d Coulomb repulsion, which manifested itself in the reduction of  $F_{dd}$ ,  $F_{pd}$ , and  $G_{pd}$  Slater integrals to 75% of their atomic values, while an unscaled value was used for the 2p core state and 3d valence state spin–orbit coupling.

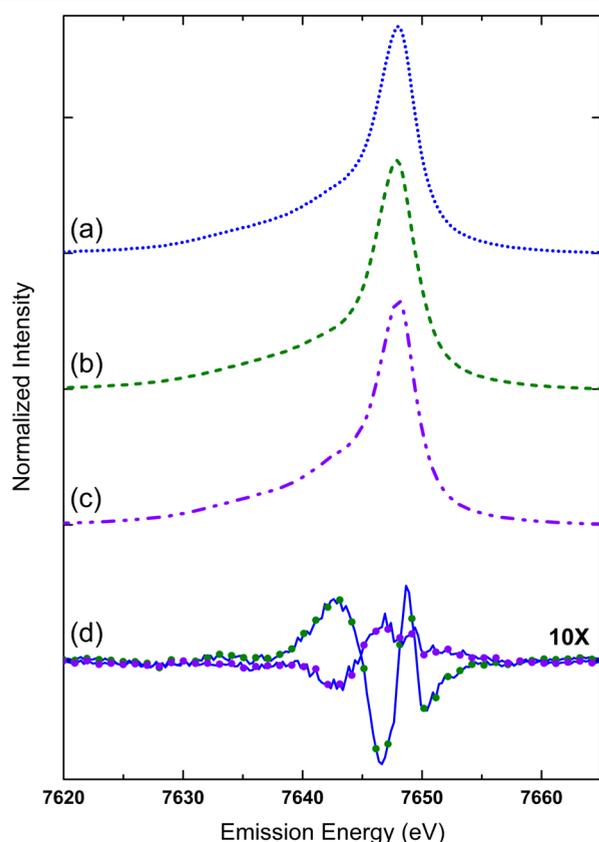
**Magnetic Measurements.** Magnetic susceptibility measurements were performed using Quantum Design’s MPMS3, a SQUID magnetometer, in an applied field of 5000 Oe and with a ramp of 5 K  $\text{min}^{-1}$  in the temperature range of 25–375 K. The Curie–Weiss law was applied in the data analysis to yield an adjusted  $\chi$  quantity, which can be related to  $\mu_{\text{eff}} = \sqrt{8\chi T}$ . When valence tautomerization is described as an equilibrium between two redox isomeric states (low- $T \rightleftharpoons$  high- $T$ ),  $\mu_{\text{eff}}^2 = f_{\text{ls}}\mu_{\text{eff,ls}}^2 + f_{\text{hs}}\mu_{\text{eff,hs}}^2$  and  $1 = f_{\text{ls}} + f_{\text{hs}}$ , where  $\mu_{\text{eff,ls}}$  is the limiting value of the effective magnetic moment for the low- $T$ , low-spin valence tautomeric state and  $\mu_{\text{eff,hs}}$  is the limiting value of the effective magnetic moment for the high- $T$ , high-spin valence tautomeric state. The mole fraction of the high- $T$ , high-spin valence tautomeric state,  $f_{\text{hs}}$ , can be determined and plotted as a function of temperature.

## RESULTS AND DISCUSSION

While different cobalt valence tautomers have been extensively investigated by various methods,<sup>12,57–59</sup> selective observations of the electronic structure at the Co center have yet to be accomplished. In this work, we utilize temperature-dependent magnetic susceptibility of the valence tautomeric complexes  $\text{Co}(\text{phen})$  and  $\text{Co}(\text{tmeda})$  as a guide for identifying their redox isomeric states and later as a complementary tool, along with our X-ray spectroscopic measurements, for understanding the Co–SQ<sup>−</sup> magnetic exchange interactions at high temperature. To begin,  $K\beta$  XES will be used to investigate and assign the oxidation and spin states of the cobalt complexes in their VT states. Followed by a combination of Co L-edge XAS and ligand-field atomic-multiplet calculations, the high- $T$  valence tautomeric state will be studied in more detail to address its electronic configuration. Last, metal–ligand interactions in valence tautomers will be discussed in the context of our element-specific X-ray- and molecular-based measurements.

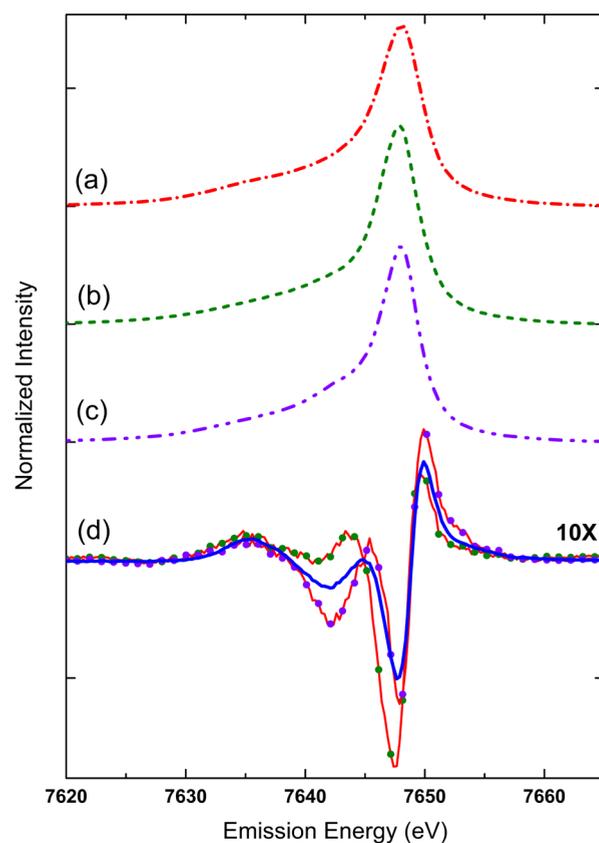
**I. Assignment of Oxidation and Spin States of Valence Tautomers at Low Temperature Using  $K\beta$  XES.** Comparing the  $K\beta$  XES spectra of  $\text{Co}(\text{phen})(3,5\text{-DBQ})_2$  and  $\text{Co}(\text{tmeda})(3,5\text{-DBQ})_2$  to those of cobalt model compounds with well-established charge and spin states provides a powerful approach to characterizing the electronic structure of the valence tautomers. We use two familiar  $\text{ls-Co}^{\text{III}}$  complexes,  $[\text{Co}(\text{en})_3]\text{Cl}_3 \cdot 2\text{H}_2\text{O}$  and  $\text{Co}(\text{acac})_3$ , for comparison because they share a similar ligand environment and coordination geometry with the cobalt valence tautomers under investigation.  $[\text{Co}(\text{en})_3]\text{Cl}_3 \cdot 2\text{H}_2\text{O}$  ( $\text{en}$  = ethylenediamine) is an excellent model complex because of its structural similarity to one of the valence tautomers,  $\text{Co}(\text{tmeda})(3,5\text{-DBQ})_2$ , whereas the acetylacetonate ( $\text{acac}$ ) ligand in  $\text{Co}(\text{acac})_3$  resembles the “ $\text{O}-\text{C}=\text{C}-\text{O}$ ” moiety in  $o$ -dioxolene.

The  $K\beta$  XES spectra of  $\text{Co}(\text{tmeda})$  and  $\text{Co}(\text{phen})$ , taken at 36 and 34 K, respectively, are compared with  $[\text{Co}(\text{en})_3]\text{Cl}_3 \cdot 2\text{H}_2\text{O}$  and  $\text{Co}(\text{acac})_3$ , as shown in Figures 2 and 3.



**Figure 2.**  $\text{Co } K\beta$  XES of (a)  $\text{Co}(\text{tmeda})(3,5\text{-DBQ})_2$  taken at 36 K (blue, dotted line), (b)  $[\text{Co}(\text{en})_3]\text{Cl}_3 \cdot 2\text{H}_2\text{O}$  (dark green, dashed line), (c)  $\text{Co}(\text{acac})_3$  (purple, dash-dot-dotted line), where parts b and c are  $\text{ls-Co}^{\text{III}}$  reference model complexes, and (d) the difference spectra, defined as  $I - I_{\text{ref}}$  of the low- $T$  valence tautomer and  $\text{ls-Co}^{\text{III}}$  model complexes:  $\text{Co}(\text{tmeda})(3,5\text{-DBQ})_2$  at 36 K with  $[\text{Co}(\text{en})_3]\text{Cl}_3 \cdot 2\text{H}_2\text{O}$  (blue solid line + dark green solid circle) and  $\text{Co}(\text{acac})_3$  (blue solid line + purple solid circle).

$\text{Co}(\text{phen})(3,5\text{-DBQ})_2$  is a classic valence tautomer studied extensively by conventional techniques such as EPR, UV–visible, Fourier transform infrared spectroscopy, and temperature-dependent magnetic susceptibility.<sup>51,54</sup> On the basis of prior studies<sup>29</sup> and magnetic susceptibility measurements in this and other work, the  $\text{Co}(\text{phen})$  system interconverts between its



**Figure 3.**  $\text{Co } K\beta$  XES of (a)  $\text{Co}(\text{phen})(3,5\text{-DBQ})_2$  taken at 34 K (red, dash-dotted line), (b)  $[\text{Co}(\text{en})_3]\text{Cl}_3 \cdot 2\text{H}_2\text{O}$  (dark green, dashed line), (c)  $\text{Co}(\text{acac})_3$  (purple, dash-dot-dotted line), where parts b and c are  $\text{ls-Co}^{\text{III}}$  reference model complexes, and (d) the difference spectra, defined as  $I - I_{\text{ref}}$  of the low- $T$  valence tautomer and  $\text{ls-Co}^{\text{III}}$  model complexes:  $\text{Co}(\text{phen})(3,5\text{-DBQ})_2$  at 34 K with  $[\text{Co}(\text{en})_3]\text{Cl}_3 \cdot 2\text{H}_2\text{O}$  (red solid line + dark green solid circle) and  $\text{Co}(\text{acac})_3$  (red solid line + purple solid circle), along with a scaled (2 $\times$ ) difference spectrum generated from the  $\text{hs-Co}^{\text{II}}$   $\text{Co}(\text{acac})_2$  and  $\text{ls-Co}^{\text{III}}$   $\text{Co}(\text{acac})_3$  complexes (blue, solid line).

two valence tautomeric states at  $T_{1/2} \approx 243$  K, whereas the  $\text{Co}(\text{tmeda})$  system transitions at  $T_{1/2} \approx 192$  K. The sharp increase in the magnetic moment of  $\text{Co}(\text{tmeda})$  around  $T_{1/2}$ , as is evident in Figure 1c, signifies a narrow temperature range over which valence tautomerization occurs. By selecting a temperature far below their  $T_{1/2}$  values, we ensure that both complexes are predominantly present in their low- $T$  valence tautomeric state,  $\text{Co}(\text{phen})(3,5\text{-DBSQ})(3,5\text{-DBCat})$  and  $\text{Co}(\text{tmeda})(3,5\text{-DBSQ})(3,5\text{-DBCat})$ , prior to X-ray exposure.

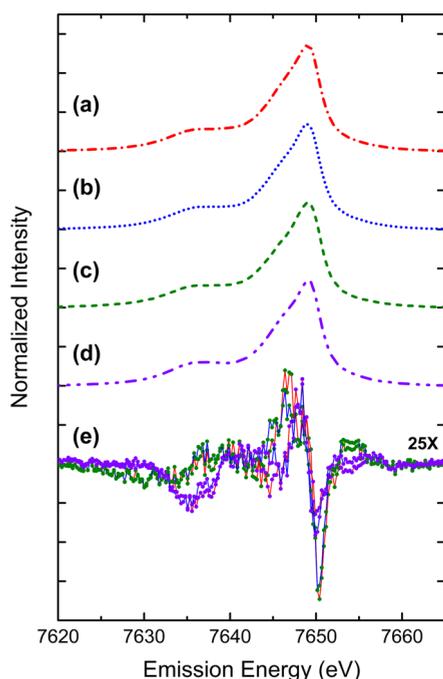
The  $K\beta$  XES spectra for  $\text{Co}(\text{tmeda})$  at 36 K,  $[\text{Co}(\text{en})_3]\text{Cl}_3 \cdot 2\text{H}_2\text{O}$  and  $\text{Co}(\text{acac})_3$ , and the differences between  $\text{Co}(\text{tmeda})$  and these model complexes can be found in Figure 2. The strong similarity between the model spectra and valence tautomer with  $\text{tmeda}$  as its coligand demonstrates that low- $T$   $\text{Co}(\text{tmeda})$  is well described by a single  $\text{ls-Co}^{\text{III}}$   $^1A_{1g}(t_{2g}^6 e_g^0)$  electronic configuration. Additionally, the shape of the difference spectra in Figure 2d does not support the presence of  $\text{hs-Co}^{\text{II}}$  and thus rules out the observation of hard X-ray-induced valence tautomerization for  $\text{Co}(\text{tmeda})$ , as discussed in more detail in the Supporting Information, section I. The equivalence of the first and second XES spectra taken at the same sample position provides further support for the absence

of hard X-ray-induced valence tautomerization for Co(tmeda) (Figure S1).

Figure 3 shows the  $K\beta$  XES spectra for Co(phen) at 34 K, [Co(en)<sub>3</sub>]Cl<sub>3</sub>·2H<sub>2</sub>O and Co(acac)<sub>3</sub>, and the differences between Co(phen) and these model complexes along with a scaled difference plot generated from the hs-Co<sup>II</sup> Co(acac)<sub>2</sub> and ls-Co<sup>III</sup> Co(acac)<sub>3</sub> complexes. Unlike Co(tmeda), the low-*T* Co(phen)  $K\beta$  XES spectrum deviates significantly from the model complex spectra, with a deviation that strongly resembles the difference spectrum generated from Co(acac)<sub>2</sub> and Co(acac)<sub>3</sub>. For this reason, we conclude that the low-*T* XES spectrum for Co(phen) shows the possibility of hard X-ray-induced VT, does not provide a clean measurement of the low-*T* valence tautomer, and cannot be used to determine the electronic structure of the low-*T* valence tautomeric state.

## II. Assignment of the Oxidation and Spin States of Valence Tautomers at High Temperature Using $K\beta$ XES.

To establish the metal spin and oxidation states in the high-*T* valence tautomers, we collected  $K\beta$  XES spectra at 375 K, a temperature well above the  $T_{1/2}$  value for both Co(phen) and Co(tmeda). Looking at Figure 4, the experimental  $K\beta$  XES



**Figure 4.** Co  $K\beta$  XES of (a) Co(phen)(3,5-DBQ)<sub>2</sub> taken at 375 K (red, dash-dotted line), (b) Co(tmeda)(3,5-DBQ)<sub>2</sub> taken at 375 K (blue, dotted line), (c) Co(phen)<sub>3</sub>·Cl<sub>2</sub>·4H<sub>2</sub>O (dark green, dashed line), (d) Co(acac)<sub>2</sub> (purple, dash-dot-dotted line), where parts c and d are hs-Co<sup>II</sup> reference model complexes, and (e) the difference spectra, defined as  $I - I_{\text{ref}}$  of the high-*T* valence tautomers and hs-Co<sup>II</sup> model complexes: Co(phen)(3,5-DBQ)<sub>2</sub> at 375 K with Co(phen)<sub>3</sub>·Cl<sub>2</sub>·4H<sub>2</sub>O (red solid line + dark green solid circle) and Co(acac)<sub>2</sub> (red solid line + purple solid circle); Co(tmeda)(3,5-DBQ)<sub>2</sub> at 375 K with Co(phen)<sub>3</sub>·Cl<sub>2</sub>·4H<sub>2</sub>O (blue solid line + dark green solid circle) and Co(acac)<sub>2</sub> (blue solid line + purple solid circle).

spectra for the high-*T* valence tautomers resemble those of Co(phen)<sub>3</sub>·Cl<sub>2</sub>·4H<sub>2</sub>O and Co(acac)<sub>2</sub>, where the latter two complexes provide well-established models for hs Co<sup>II</sup> with similar coordination geometry and bonding. In the hs-Co<sup>II</sup> XES spectra, the  $K\beta_{1,3}$  main peak has an asymmetric line shape with more spectral intensity at energies lower than the peak at

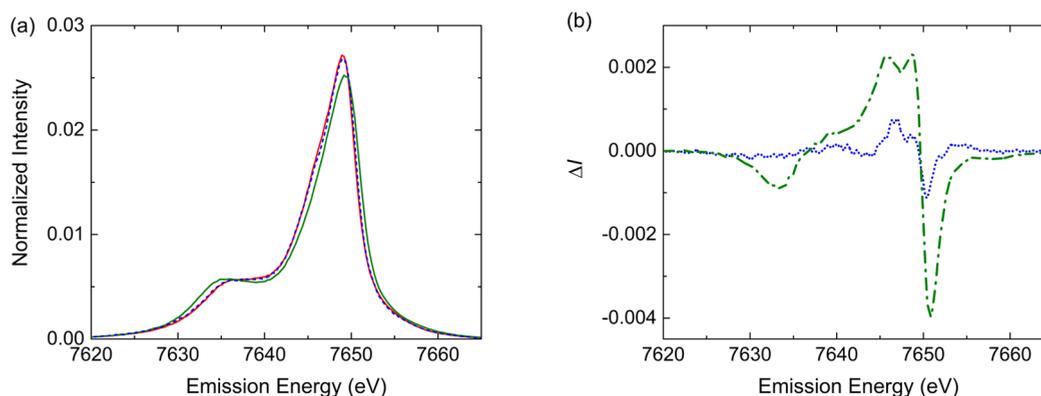
7649.0 eV. In addition, the presence of a broad  $K\beta'$  shoulder at 7636.0 eV provides the expected spectral signature associated with a hs-Co<sup>II</sup> configuration, which is absent for a ls-Co<sup>II</sup> XES (Figure S11).<sup>60</sup> The two hs-Co<sup>II</sup> model complexes describe the spectral line shape of both high-*T* valence tautomers equally well, as is evident in their difference spectra in Figure 4e. Taken together, the high-*T* Co(phen) and Co(tmeda) valence tautomers are positively identified to be in a hs  $d^7$  configuration,  $t_{2g}^5 e_g^2$ , a conclusion supported by X-ray crystallography and other techniques.<sup>51</sup>

Results very similar to those observed in Co(phen) and Co(tmeda) have also been obtained for other cobalt valence tautomers in the high-*T* state, including Co(bpy), Co(NO<sub>2</sub>-phen), Co(tmmda), and Co(tmpda), where bpy = 2,2'-bipyridine, NO<sub>2</sub>-phen = 5-nitro-1,10-phenanthroline, tmmda = *N,N,N',N'*-tetramethylmethylenediamine, and tmpda = *N,N,N',N'*-tetramethylpropylenediamine (Figure S6). In all cases, the valence tautomers share the same spectral features and line shapes as the hs-Co<sup>II</sup> model complexes, despite the varying  $\sigma$ -donor properties of their distinct nitrogen coligands. On the basis of the cobalt-specific XES data in this work, the  $\sigma$ -donor abilities of the nitrogen coligands do not tip the thermodynamic balance in favor of a ls-Co<sup>II</sup> form at high temperature, as previously suggested by Adams et al.<sup>61</sup> The influence of the coligand on the transition temperature<sup>51,54,62</sup> does not lead to a change in the electron correlation at the Co center because the high-*T* valence tautomers all have the same hs-Co<sup>II</sup> electronic configuration. Last, temperature-dependent  $K\beta$  XES measurements show two-state thermodynamics, as expected (see Figure S4).

## III. Absence of hs-Co<sup>III</sup> Electron Configurations in the High-*T* Valence Tautomeric State.

Attempts to describe the electronic structure of cobalt valence tautomers have highlighted the potential importance of strong configuration interaction. In a combined theoretical and spectroscopic study by LaBute et al.,<sup>39</sup> they concluded that both hs-Co<sup>II</sup> and hs-Co<sup>III</sup> configurations contribute to the electronic structure of the high-*T* valence tautomers. To explore this possibility further, the  $K\beta$  XES spectra of Co(phen)<sub>3</sub>·Cl<sub>2</sub>·4H<sub>2</sub>O, one of the hs-Co<sup>II</sup> model complexes, and CoF<sub>3</sub>, a rare example of a hs-Co<sup>III</sup> octahedral species, are shown in Figure 5a. Compared to Co(phen)<sub>3</sub>·Cl<sub>2</sub>·4H<sub>2</sub>O, the positions of the  $K\beta'$  shoulder for CoF<sub>3</sub> (7635.0 vs 7636.0 eV for the former) and the  $K\beta_{1,3}$  main peak (7649.4 vs 7649.0 eV) are shifted to lower and higher emission energies, respectively, relative to the spectral center. Despite their dissimilar metal–ligand bonding interaction, the aforementioned spectral difference between CoF<sub>3</sub> and Co(phen)<sub>3</sub>·Cl<sub>2</sub>·4H<sub>2</sub>O is confirmed by, and consistent with, our calculated hs-Co<sup>III</sup> and hs-Co<sup>II</sup> XES spectra (Figure S11). The spectral differences between Co(phen) at high-*T* and the two hs-Co model spectra appear in Figure 5b. On the basis of their variational configuration interaction analysis of the Co K preedge XAS, LaBute et al.<sup>39</sup> concluded that Co(phen) at high-*T* is comprised of 30% hs Co<sup>III</sup> and 70% hs Co<sup>II</sup>. The X-ray emission spectra of molecules with multiple configurations can be modeled as a weighted sum of single-electron configuration spectra, with amplitudes ( $\alpha$ 's) proportional to the contribution of each configuration to the overall electronic structure.<sup>39,63,64</sup> A multiconfigurational representation of the high-*T* valence tautomeric state is as follows:

$$|\psi_{\text{high-}T \text{ VT}}\rangle = \alpha_1 |\text{hs Co}^{\text{III}}, d^6\rangle + \alpha_2 |\text{hs Co}^{\text{II}}, d^7\rangle$$

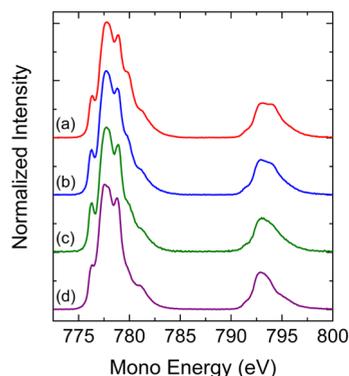


**Figure 5.** (a) Co  $K\beta$  XES of  $\text{CoF}_3$  (dark green, solid line), a  $\text{hs-Co}^{\text{III}}$  reference model,  $\text{Co(phen)}_3\cdot\text{Cl}_2\cdot 4\text{H}_2\text{O}$  (blue, dashed line), a  $\text{hs-Co}^{\text{II}}$  reference model, and  $\text{Co(phen)}_3(3,5\text{-DBSQ})_2$  taken at 375 K (red, solid line). (b) Difference spectra of  $\text{Co(phen)}_3(3,5\text{-DBSQ})_2$  with  $\text{CoF}_3$  (dark green, dash-dotted line) and  $\text{Co(phen)}_3\cdot\text{Cl}_2\cdot 4\text{H}_2\text{O}$  (blue, dotted line). Note that  $\Delta I$  is defined as  $I - I_{\text{ref}}$ .

where  $|\psi_{\text{high-}T\text{VT}}\rangle$  is the complete wave function of the high- $T$  valence tautomeric state and  $\alpha_1$  and  $\alpha_2$  are the coefficients for the corresponding electron configurations  $|\text{hs Co}^{\text{III}}, d^6\rangle$  and  $|\text{hs Co}^{\text{II}}, d^7\rangle$ .

To test the significance of a  $\text{hs-Co}^{\text{III}}$  configuration, we fit the  $\text{Co(phen)}$   $K\beta$  XES data to a linear sum of the  $\text{Co(phen)}_3\cdot\text{Cl}_2\cdot 4\text{H}_2\text{O}$  and  $\text{CoF}_3$  spectra. Zero spectral weight for the  $\text{hs-Co}^{\text{III}}$  spectrum yields the minimum  $\chi^2$  fit and provides no evidence for a  $\text{hs-Co}^{\text{III}}$  configuration in the high- $T$  valence tautomer.

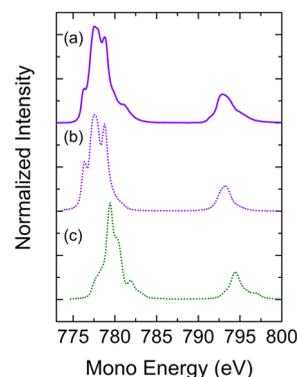
**IV. Co L-Edge XAS of the High- $T$  Valence Tautomeric State.** To further test if a  $\text{hs-Co}^{\text{II}}$  electron configuration alone best describes the metal center in valence tautomers, we used Co L-edge XAS to investigate both  $\text{Co(phen)}$  and  $\text{Co(tmeda)}$  at high temperature. Unlike  $K\beta$  XES, which probes the 3d valence electrons via its sensitivity to the 3p3d Coulomb and exchange interaction in the final states, Co L-edge XAS provides complementary information by examining the 3d unoccupied levels via optical excitation of a 2p core electron directly into 3d unfilled orbitals in an electric dipole-allowed transition.<sup>65–69</sup> Figure 6 shows the Co L-edge XAS of  $\text{Co(phen)}$  and  $\text{Co(tmeda)}$  at room temperature, as well as the  $\text{hs-Co}^{\text{II}}$  model complexes:  $\text{Co(phen)}_3\cdot\text{Cl}_2\cdot 4\text{H}_2\text{O}$  and  $\text{Co}_4(3,5\text{-DBSQ})_8$ . The spectra of these complexes all show strong similarities. The additional intensity at 779.8 eV for  $\text{Co(phen)}$  can be explained by a small population of the low- $T$



**Figure 6.** Co L-edge XAS of (a)  $\text{Co(phen)}_3(3,5\text{-DBSQ})_2$  (red, solid line), (b)  $\text{Co(tmeda)}_3(3,5\text{-DBSQ})_2$  (blue, solid line), (c)  $\text{Co(phen)}_3\cdot\text{Cl}_2\cdot 4\text{H}_2\text{O}$  (dark green, solid line), and (d)  $\text{Co}_4(3,5\text{-DBSQ})_8$  (purple, solid line). Both parts c and d are  $\text{hs-Co}^{\text{II}}$  reference model complexes. All of the spectra were taken at room temperature.

valence tautomers still present at 298.0 K, given its slightly higher  $T_{1/2}$  value than that of  $\text{Co(tmeda)}$ . The fast rate of soft X-ray-induced valence tautomerization precluded the use of Co L-edge XAS to characterize the valence tautomeric electronic structure at low- $T$ .

We also used ligand-field atomic-multiplet calculations to support the conclusions drawn from our X-ray measurements. The Co L-edge XAS of  $\text{Co}_4(3,5\text{-DBSQ})_8$  is compared side by side with the calculated  $\text{hs-Co}^{\text{II}}$  and  $\text{hs-Co}^{\text{III}}$  spectra (Figure 7).



**Figure 7.** Comparison of Co L-edge XAS of (a)  $\text{Co}_4(3,5\text{-DBSQ})_8$  taken at room temperature (purple, solid line), a  $\text{hs-Co}^{\text{II}}$  reference compound, (b) calculated  $\text{hs Co}^{\text{II}}$  (purple, dotted line), and (c) calculated  $\text{hs Co}^{\text{III}}$  (dark green, dotted line). Note that the calculated spectra in parts b and c were scaled to the maximum peak height of the experimental data and red-shifted by 1.32 eV for comparison.

The calculations used an  $O_h$  molecular symmetry and a  $10 Dq = 1.0$  eV. The spectral shape for the calculated  $\text{hs Co}^{\text{II}}$  compares well with our L-edge spectra, but this is not the case with the calculated  $\text{hs-Co}^{\text{III}}$  spectrum. Note that the simulated  $K\beta$  XES spectra for  $\text{hs Co}^{\text{II}}$  and  $\text{hs Co}^{\text{III}}$  using the same atomic parameter settings compare well with their respective experimental data (Figure S11). Because an appropriate  $\text{hs-Co}^{\text{III}}$  model complex is lacking for the Co L-edge XAS measurements, we instead carried out extensive ligand-field atomic-multiplet calculations with different  $10 Dq$  settings and Slater integral values (Figures S7–S9) to derive a comprehensive overview of the expected spectral shapes and features. These calculations further support the assignment of the high- $T$  valence tautomer's electronic structure to a single determinant  $\text{hs Co}^{\text{II}}$ .

**V. Magnetic Exchange Interaction between the Co Center and Semiquinone Ligands in Valence Tautomers.** Using  $K\beta$  XES, the spin state of the Co center in the two valence tautomeric states can be directly and unambiguously assigned. On the basis of the observation that the low- $T$  state exists as a  $d^6$ ,  $ls\text{-Co}^{\text{III}}$  center and the high- $T$  state as a  $d^7$ ,  $hs\text{-Co}^{\text{II}}$  center, there are nominally zero and three unpaired electrons found at the metal site, respectively. By relating the electronic information at the Co center determined using X-ray spectroscopies with the molecular magnetic moments of valence tautomers measured with SQUID magnetometry, new insight pertaining to the magnetic exchange interactions between the unpaired electrons on the metal center and those on the ligands can be deduced. The magnetic plots of  $\text{Co}(\text{phen})$  and  $\text{Co}(\text{tmeda})$  show simple two-state thermodynamics with weak intermolecular magnetic exchange interactions in the temperature range that we have studied with X-ray spectroscopy, unlike the strong temperature dependence of the magnetic properties that has been seen in similar  $\text{Co}(\text{SQ})_2$  valence tautomers at low temperature.<sup>70,71</sup> Taken together, the magnetic behaviors of our system allow for a straightforward interpretation in the temperature range of 25–375 K.

As observed in previous studies, the low- $T$  valence tautomeric complexes have magnetic moments and EPR values consistent with those of one unpaired electron.<sup>72,73</sup> The  $K\beta$  emission spectra at low- $T$  have been assigned to a  $ls\text{-Co}^{\text{III}}$  single electronic configuration, confirming that the unpaired electron is localized on a singly reduced semiquinone ligand.<sup>61,72,74,75</sup>

The  $K\beta$  emission spectra at high- $T$  have been assigned to a  $hs\text{-Co}^{\text{II}}$  quartet electronic structure, with its three unpaired electrons distributed between the  $t_{2g}$  ( $d\pi$ ) and  $e_g$  ( $d\sigma$ ) orbitals. It is coordinated to two semiquinone ligands, each of which carries one unpaired electron in its valence orbital. Consequently, a total of five unpaired electrons are available to interact in the complex at high temperature. Consistent with previous magnetometry measurements,<sup>51</sup> the molecular spin moment is best assigned to  $S = 3/2$ <sup>76–80</sup> and thus is inconsistent with the  $S = 1/2$  expected for strong antiferromagnetic coupling or the  $S = 5/2$  expected for strong ferromagnetic coupling between the  $hs$  quartet  $d^7$  metal and doublet ligands. Given the absence of  $ls\text{-Co}^{\text{II}}$  configurations at high temperature, we can also rule out the possibility of a ferromagnetic molecular quartet with one unpaired electron on cobalt and each semiquinone.

In order to explain the metal–ligand magnetic exchange interaction based on a  $hs\text{-Co}^{\text{II}}$  center in valence tautomers at high- $T$ , two alternate scenarios are possible. In the first and most probable scenario, the magnetic exchange interaction between cobalt and semiquinones is weak, on the order of a few millielectronvolts, such that the unpaired electrons on the ligands are weakly coupled to and oriented randomly with respect to the cobalt spin vector, an idea that was proposed by Pierpont and co-workers.<sup>81–84</sup> With a small-energy-level spacing between the  $S = 1/2, 3/2, 5/2$  electronic states, all spin configurations will be thermally populated at room temperature. Alternatively, the high- $T$  state can be described as having a strong antiferromagnetic interaction between the spins on the semiquinone ligands. In this case, the two unpaired electrons on the ligands are presumed to pair opposite to each other, giving the ligands a total spin of  $S = 0$  and yielding a coordination complex with a total spin of  $S = 3/2$ . As a general conclusion from studies of  $\text{M}(\text{SQ})_3$  complexes, where the metal ion (M) is paramagnetic,  $\text{M-SQ}^-$  coupling is strong and

antiferromagnetic for unpaired  $d\pi$  metal electrons.<sup>11,12,85</sup> However, the  $hs$   $d^7$  center in our case, with only one unpaired spin in its  $d\pi$  orbitals, does not support the case for strong  $\text{M-L}$  interactions.<sup>86,87</sup> Additionally, strong metal-mediated ligand–ligand interactions happen more often when the metal center is diamagnetic, not paramagnetic.<sup>88,89</sup> This reflects the shorter metal–ligand bond lengths in low-spin complexes, allowing for the proper orbital alignment to take place.<sup>90</sup> Again, that is not the case with the  $hs\text{-Co}^{\text{II}}$  center in the high- $T$  valence tautomeric state; with its  $d\sigma$  orbitals half-filled, it is expected that a longer  $\text{Co-O}$  bond distance would result in weak  $\text{Co-SQ}^-$  bonding interactions, which matches with previous X-ray crystallography of the same valence tautomers studied at the two extreme temperatures.<sup>51</sup> Consequently, a scenario of strong antiferromagnetic interaction between two spatially separated  $\text{SQ}^-$ , mediated by a metal center with which they may only be weakly interacting, does not seem chemically or electronically likely. Nonetheless, this has been a popular theory among many groups.<sup>20,88,91</sup> Given our robust determination of the cobalt spin and oxidation states, in combination with experimental precedents on strong  $\text{M-SQ}^-$  interactions in  $\text{M}(\text{SQ})_3$  complexes and metal-mediated ligand–ligand interactions, we conclude that weak magnetic exchange interactions<sup>92</sup> between the Co and  $\text{SQ}^-$  ligands in the high- $T$  valence tautomeric state provide the most probable description of the exchange interaction in these complexes.

## ■ CLOSING REMARKS

On the basis of the cobalt-specific, spin-sensitive  $K\beta$  XES and L-edge XAS, in conjunction with ligand-field atomic-multiplet calculations, it is evident that the Co center in both valence tautomeric states can be accurately modeled with single electronic configurations,  $ls$   $\text{Co}^{\text{III}}$  in the low- $T$  state and  $hs$   $\text{Co}^{\text{II}}$  in the high- $T$  state. We find no evidence for either  $ls\text{-Co}^{\text{II}}$  or  $hs\text{-Co}^{\text{III}}$  electronic configurations, as previously postulated to be of importance by Adams et al.<sup>61</sup> and LaBute et al.<sup>39</sup> This conclusion applies to numerous cobalt valence tautomers with chemically distinct N-donor coligands. For a limited number of nitrogen coligands, we observe that the coligand predominantly influences the transition temperatures, rather than the general electronic features, of the valence tautomers. In order to better understand the contribution and influence of nitrogen coligands in VT, further systematic investigations of different electron-withdrawing and -donating groups are needed.

By extending  $K\beta$  XES to the time domain, one can envision future studies dissecting the mechanism of photoinduced VT.<sup>93–96</sup> To convert the complex from the  $ls\text{-Co}^{\text{III}}$  state coordinated to a catecholate and a semiquinone ligand to the  $hs\text{-Co}^{\text{II}}$  state coordinated to two semiquinone ligands, one electron from the  $\text{Cat}^{2-}$  ligand has to be transferred to the metal via a ligand-to-metal charge-transfer excitation, thereby reducing the charge at the Co center and involving either a simultaneous or a sequential spin transition. While many research laboratories have attempted to address the thermally driven phenomenon via theoretical calculations,<sup>39,40,61</sup> this paper opens up the possibility of future photochemical studies—as previously demonstrated in various SCO systems<sup>97–101</sup>—on these transient spin and charge states of valence tautomers using  $K\beta$  XES in a time-resolved manner.

## ■ ASSOCIATED CONTENT

### ● Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acs.inorgchem.6b01666.

Discussion of X-ray-induced VT, additional  $K\beta$  XES and difference emission spectra demonstrating the thermally driven valence tautomeric interconversion, and detailed ligand-field atomic-multiplet calculations of a  $hs\text{-Co}^{\text{III}}$  and a  $ls\text{-Co}^{\text{II}}$  center, along with concise descriptions of these spectral and calculated data (PDF)

## ■ AUTHOR INFORMATION

### Corresponding Authors

\*E-mail: wliang6@stanford.edu.

\*E-mail: kgaffney@slac.stanford.edu.

### ORCID

H. Winnie Liang: 0000-0001-6548-2061

### Notes

The authors declare no competing financial interest.

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