Pressure-induced spin reorientation and spin state transition in SrCoO₃


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A series of high-pressure measurements including resistivity, ac and dc magnetic susceptibility, neutron powder diffraction, and synchrotron x-ray diffraction were performed to investigate the crystal structure and electronic states of the itinerant ferromagnet SrCoO₃. Two pressure-induced phase transitions were observed at about 1.1 and 45 GPa, corresponding to a spin reorientation and a spin-state transition, respectively, while the cubic crystal structure was stable with pressure up to 60.0 GPa. The origins of these electronic state changes are discussed and rationalized in the light of first-principles calculations.

DOI: 10.1103/PhysRevB.92.195147 PACS number(s): 61.50.Ks, 71.15.Mb, 75.25.—j, 75.50.—y

I. INTRODUCTION

Cobalt-based oxides continue to attract great attention due to their versatile charge, spin, and orbital degrees of freedom as well as the associated fascinating physical properties such as giant magnetoresistance [1–3], high-performance thermoelectricity [4], unconventional superconductivity [5], and large linear magnetoelectric effect [6]. As is well known, the spin states of Co ions are strongly dependent on the subtle balance between the crystal field energy and the Hund’s rule exchange. As a consequence, various spin states are observed in cobalt oxides [7–15]. Since external stimuli can significantly change such an energy balance, Co ions in real materials can display a series of interesting transitions between different spin states under different temperature and pressure conditions. A well-known example is the low-spin (LS) to high-spin (HS) transition via a possible intermediate-spin (IS) state that is observed in the rhombohedral perovskite LaCoO₃ with increasing temperature [16–21]. In addition, Pr₀.₅Ca₀.₅CoO₃ exhibits a pressure-induced IS to LS transition, leading to a coherent change between a magnetic state with metallic conduction and a paramagnetic insulating state [22]. Moreover, pressure-induced crystal structural and spin-state phase transitions have also been found in other Co oxides such as BiCoO₂ and SrRu₀.₅Co₀.₅O₃, etc. [23,24].

In Co-based perovskite oxides, SrCoO₃ stands out due to its undistorted cubic perovskite structure with space group Pm₃m as well as the strong ferromagnetic (FM) interaction with a high Curie temperature (Tᵣ) above room temperature (∼305 K measured for a single crystal in a field of 0.5 T) [12,13]. Although the spin state of Co⁴⁺ in SrCoO₃ has been a subject of active debate, recent experimental studies on high-quality single crystals provide convincing evidence for the IS state [12], in accordance with atomic multiplet theoretical calculations based on the Co-2p x-ray absorption spectrum [25,26]. Moreover, theoretical analysis has revealed that (1) there exist strong p-d hybridizations and oxygen-hole effects due to the large negative charge-transfer energy between O-2p and Co-3d orbitals; (2) the electronic configuration of the Co⁴⁺ ion with LS state (i.e., the ground state consisting of 8% d⁵, 67% d⁶L, and 25% d⁷L²) where the L represents an oxygen hole; (3) the FM metallic state of SrCoO₃ mainly originates from the strong coupling between the itinerant oxygen holes and the high-spin Co-3d⁶ ions (i.e., the dominant configuration d⁶L). The presence of the IS state in SrCoO₃ is thus indicative of the competing crystal field and Hund’s-rule interactions, the relative strength of which can be tuned by external perturbations. It is therefore interesting to study pressure effects on the magnetism and spin state of SrCoO₃. In this paper, we describe a wide variety of high-pressure measurements such as ac and dc magnetic susceptibility, electric resistivity, neutron powder diffraction (NPD) and synchrotron x-ray diffraction (SXRD) as well as theoretical calculations performed to characterize the structure and electronic properties of SrCoO₃ under high pressure. Two pressure-induced magnetic transitions, namely, a spin reorientation and a spin state transition were observed at different pressures.

II. EXPERIMENTAL AND CALCULATION DETAILS

High-quality SrCoO₃ was prepared by a two-step method as described in detail in Ref. [12]. Oxygen content was found to be almost stoichiometric (2.95 ± 0.02) by a thermogravimetric
analysis [12]. Pressure-dependent resistivity ($\rho$) and ac magnetic susceptibility ($\chi_{ac}$) were measured using a cubic-anvil-type high-pressure apparatus at room temperature and under zero magnetic field. The temperature- and field-dependent dc magnetic susceptibility ($\chi_{dc}$) measurements were carried out on a commercial superconducting quantum interference device magnetometer (Quantum Design) at various pressures. The pressures were generated by a piston-cylinder high-pressure system with glycerin as pressure-transmitting medium. The pressure values were calibrated based on the superconducting transition temperature of metallic Pb.

The high-pressure ($P < 5.4$ GPa) NPD measurements were carried out using a Paris-Edinburgh cell installed at the HRPT high-resolution diffractometer in the Swiss Spallation Neutron Source (SINQ), Paul Scherrer Institute [27]. The measurements were conducted at 200 K, a temperature at which the pressure can still be considered as hydrostatic, and using a neutron wavelength $\lambda = 1.4931$ Å. NaCl was mixed in with the powder sample and the observed pressure dependence of its cubic lattice parameter was used to calibrate the applied sample pressure. A 4:1 mixture of deuterated methanol-ethanol was used as pressure-transmitting medium. The NPD data were analyzed using the Rietveld package FULLPROF SUITE [28]. The primary room-temperature high-pressure SXRD characterization was performed at the Beijing Synchrotron Radiation Facility, followed by detailed measurements under high pressure conducted at the beamline 16BM-D in the Advanced Photon Source, Argonne National Laboratory. The wavelength used for SXRD measurements was $\lambda = 0.3099$ Å. Pressure was generated by a diamond anvil cell with a pair of 300 μm culet anvils. Neon gas was used as pressure-transmitting medium and the pressure values were calculated by the ruby fluorescence method. The SXRD data were analyzed using the Rietveld refinement program GSAS [29].

First-principles calculations based on density functional theory (DFT) were performed for SrCoO$_3$. We used the generalized gradient approximation and the projected augmented wave method with a plane wave basis set as implemented in the Vienna ab initio Simulation Package [30–32]. A kinetic energy cut off was set to be 500 eV for the plane wave basis and the structural optimization was sampled with a 16 $\times$ 16 Monkhorst-Pack k-point grid in combination with the tetrahedron method. All coordinates of atomic positions of SrCoO$_3$ were fully relaxed until the forces became small than 0.01 eV Å$^{-1}$. The convergence of our calculations was checked carefully. The Wannier projection on DFT-calculated Co-d Bloch waves was performed with the VASP2WANNIER90 interface employing WANNIER90 for constructing maximally localized Wannier orbitals. Through the Wannier projection, we obtained all hopping terms and crystal field energies, and constructed a tight-binding Hamiltonian, which exactly reproduced the DFT-calculated Co-d bands. Magnetic anisotropy energies under pressures were calculated by noncollinear DFT, where spin-orbit coupling interactions were taken into account.

### III. Results and Discussion

Figure 1(a) shows the pressure dependence of the resistivity and the ac magnetic susceptibility of SrCoO$_3$ measured at zero field and room temperature. With increasing pressure, both the $\rho$ and $\chi_{ac}$ exhibit an anomaly at a critical pressure ($P_{C1}$) about 1.1 GPa, indicating the occurrence of a pressure-induced variation in the magnetic and/or electronic properties (as shown later, the crystal structure remains unchanged). To further characterize this phase transition, dc magnetization was measured at different pressures. Based on the temperature-dependent magnetization measured at 0.1 T shown in Fig. 1(b), we find that the FM Curie temperature almost linearly increases with pressure at a rate of $\sim$13 K/GPa (see the inset), revealing that pressure enhances the FM double-exchange interaction occurring between Co-3d$^6$ electrons and oxygen holes. Note that the $T_c$ value observed here at ambient
pressure is slightly lower than that reported in Ref. [12]. This is partly due to the difference in the magnetic field value (0.5 and 0.1 T) used for the estimate of $T_C$. In addition, the temperature-dependent magnetization exhibits different features below and above $P_{C1}$. When pressure is applied below $P_{C1}$ (e.g., at 0.48 and 0.82 GPa), the magnetization shows a smooth temperature dependence below $T_C$. In sharp contrast, when the applied pressure is close to or larger than $P_{C1}$ (e.g., at 1.04 and 1.28 GPa), the magnetization experiences a clear drop at a temperature $T_{SR}$ (the spin reorientation temperature) which increases with pressure. At the first glance, one may attribute this variation to a possible spin-state transition as observed in many other cobalt oxides. However, if we check the field-dependent magnetization, the saturated magnetic moments obtained both below and above $P_{C1}$ are almost identical [see the inset of Fig. 1(c)], in contrast with the behavior expected for the occurrence of a spin-state transition at this critical pressure. Furthermore, above $P_{C1}$, when the low-field magnetization is examined in detail, one can find an S-shaped metamagnetic behavior below $T_{SR}$, whereas canonical magnetization curve is observed above $T_{SR}$, as represented in Fig. 1(c). These observations could suggest that the easy magnetization axis of SrCoO$_3$ changes under high pressure, giving rise to a spin-reorientation transition as will be demonstrated by our theoretical calculations.

In the case of a spin-state transition, changes in both the Co-O interatomic distances and the value of the Co magnetic moment are expected. Neutron powder diffraction is a desirable technique for this kind of investigation because it can obtain these parameters simultaneously. We therefore performed high-pressure NPD at 200 K (<$T_{SR}$) at pressures both below and above $P_{C1}$ to get additional insight concerning the origin of the pressure-induced phase transition observed around 1.1 GPa. As shown in Fig. 1(b), the Co$^{4+}$ moment is nearly saturated at 200 K and the saturation will most likely be maintained under the application of pressure, taking into account the linear increase of Curie temperature $T_C$ with $P$. Since the intensity of magnetic reflections is proportional to the square of the ordered Co$^{4+}$ magnetic moment, any change in moment will be easier to detect at 200 K than at room temperature, where the ordered magnetic moment is about 1/3 of the saturation value. Figures 2(a) and 2(b) show the Rietveld refinements of the NPD patterns obtained at ambient pressure (<$P_{C1}$) and 5.4 GPa (> $P_{C1}$), respectively. The pressure dependence of the unit cell volume obtained from such refinements is shown in Fig. 2(c). In other cobalt oxides such as BiCoO$_3$ [23] or GdBaCo$_2$O$_{5.5}$ [33], a spin-state transition from a higher- to a lower-spin state has been shown to lead to a discontinuous shrinking in unit cell volume. In the present case, however, the unit cell volume smoothly decreases with increasing pressure, and no obvious anomaly is observed up to the highest pressure value 5.4 GPa used in this experiment. The volume-pressure relationship can be well fitted based on the Birch-Murnaghan equation of state with $P = 1.5B_0(x^{-7/3} - x^{-5/3})(1 + 0.75(B'/4)(x^{-2/3} - 1))$ [see Fig. 2(c)]. Here $x = V/V_0$ is the normalized volume, $V_0$ is the unit cell volume at pressure $P = 0$ GPa, and $B_0$ and $B'$ are the bulk modulus and its derivative with respect to the pressure. The fitted value of $B_0$ is 144 ± 2 GPa with fixed $B' = 4$. A further argument supporting the absence of a spin-state transition at 1.1 GPa is that the magnetic contribution to the Bragg reflections is found to be nearly unchanged with pressure up to 5.4 GPa. As a consequence, the refined value of the Co$^{4+}$ magnetic moment shown in Fig. 2(c) takes a nearly constant value very close to ~1.7$\mu_B$, compatible with both a highly hybridized IS spin state and the results of the magnetization measurements [see Fig. 1(b)] [12]. Our high-pressure NPD results thus solidly confirm both the stability of the crystal structure and the spin state of SrCoO$_3$ below 5.4 GPa.
configuration. Clearly, the pressure can enhance 36.8 GPa and above 46.0 GPa, respectively.

cell volume. The red and blue solid lines are the fitting results to the allowed Bragg reflections (ticks). (c) Pressure dependence of the unit line, and difference (blue line) profiles are shown together with the observed (black crosses), calculated (red medium neon. (b) SXRD pattern and Rietveld refinement result at a typical pressure 46.0 GPa. Observed (black crosses), calculated (red

collection at various pressures and room temperature. The asterisks stand for sample pressures up to 60.0 GPa. Figure 3(a) shows some typical SXRD patterns collected at selected pressures. As the pressure increases, the peak positions systematically shift to higher diffraction angles due to the volume compression, and no structural phase transformation is observed even when the pressure is increased to 60.0 GPa. Figure 3(b) shows the Rietveld structure refinement at a representative pressure of 46.0 GPa, at which the SXRD pattern can still be well fitted on the basis of a simple cubic perovskite structure with space group \( Pm\bar{3}m \). However, when the pressure dependence of the unit cell volume is examined as shown in Fig. 3(c), a clear anomaly is observed at another critical pressure \( P_{C2} \) around 45 GPa. The change of spin state in other cobalt perovskite oxides is accompanied by an anomaly in unit cell volume and/or Co-O interatomic distances due to the different ionic radii at different spin states even in the case of no change of crystal symmetry [21,24,34]. Thus, the anomaly at \( P_{C2} \) suggests the occurrence of a spin-state variation from IS to LS in SrCoO₃. According to the Birch-Murnaghan equation of state mentioned above, the room-temperature bulk modulus below \( P_{C2} \) is fitted to be 164 ± 4 GPa, which is comparable to those observed in most perovskite oxides [35–37]. Above \( P_{C2} \), however, the bulk modulus \( B_0 \) considerably increases to 229 ± 19 GPa. In other cobalt perovskite oxides, the pressure-induced spin-state transition also gives rise to a sharp increase in \( B_0 \) [36]. In comparison, the anomaly in unit cell volume and the significant increase of \( B_0 \) in the present SrCoO₃ sample may suggest the occurrence of a pressure-induced spin-state transition from IS to LS. Note that since the unit cell volume of SrCoO₃ does not show a first-order-like drop with pressure across \( P_{C2} \), the associated IS-LS transition is more likely continuous over a wider pressure region (36.8–46.0 GPa) as shown in Fig. 3(c).

To get deeper insight into the pressure effects on the magnetic state in the present cubic perovskite SrCoO₃, first-principles calculations based on density functional theory were performed. Figure 4(a) shows the pressure-dependent crystal field energy (10\( Dq \)) and the \( p-d \) transfer interaction (|\( t | \)) between Co-3\( d \) and O-2\( p \) orbitals. Obviously, both of them increase with increasing pressure. As mentioned above, the FM interaction of SrCoO₃ mainly originates from the double exchange between the Co-3\( d^6 \) electrons and oxygen holes in the dominant \( d_{xy} \) configuration. Clearly, the pressure can enhance the \( p-d \) transfer interaction significantly [Fig. 4(a)]. As a result, the FM Curie temperature of SrCoO₃ increases with increasing pressure, consistent with experimental observation [Fig. 1(b)]. On the other hand, the crystal field energy also increases with elevating pressure, which can result in a spin-state transition if the pressure is high enough. Actually, when we check the total system energy as a function of pressure for the IS and LS states, respectively, we can find a crossover around 50 GPa where IS and LS states are strongly mixed and the transition of spin state from IS to LS is completely achieved by using higher pressure as shown in Fig. 4(b). This supports the pressure-induced and continuous-like spin-state transformation as suggested by the SXRD experiments [Fig. 3(c)]. Also we calculated the pressure dependence of the unit cell volume, and found that the bulk modulus significantly increases from \( B_0 = 166 ± 2 \) GPa to \( B_0 = 194 ± 2 \) GPa across \( P = P_{C2} \). This result is qualitatively in accordance with the experimental result [Fig. 3(c)].

Furthermore, we calculated the energy difference with the easy magnetization axis along the [110] and [111] directions...
relative to the [100] direction under high pressure. As shown in Fig. 4(c), the [111] direction is favorable in energy at lower pressures. However, when pressure increases to about 3.7 GPa, the [100] direction becomes more stable in energy. Therefore, the anomalous variations observed in resistivity and ac and dc magnetic susceptibility curves near $P_{C1}$ (see Fig. 1) may be attributed to a pressure-induced spin-reorientation transition.

IV. CONCLUSIONS

In summary, pressure effects on the crystal structure and the Co$^{4+}$ spin state were studied in detail for SrCoO$_3$ single crystals using different experimental techniques and first-principles calculations. The FM Curie temperature of SrCoO$_3$ increases with pressure due to an enhancement of the FM double-exchange interaction between Co-3$d^6$ electrons and oxygen holes. Although the cubic perovskite structure of SrCoO$_3$ is stable with pressure up to 60.0 GPa, two pressure-induced changes in the electronic system are observed at 1.1 and 45 GPa. The first of them is attributed to a spin reorientation due to a pressure-driven change of the easy magnetization axis, whereas the second one is most likely a pressure-induced spin-state transition. In contrast with the first-order spin-state transitions observed in most cobalt oxides, the present one appears to be continuous, and to occur over a wide pressure region.

ACKNOWLEDGMENTS

This work was supported by the Key Project of Chinese National Programs for Fundamental Research and Development (973 program) (Grants No. 2014CB921500 and 2015CB921303), and the Strategic Priority Research Program of the Chinese Academy of Sciences (Grants No. XDB07030300 and XDB07020200). W.Y. acknowledges the support by the NSAF (Grant No. U1530402). Y.-W.L. and Q.Z. acknowledge the support by the NSFC (Grants No. 11574378, No. 11474246, and No. 11204265), Foundation of the Jiangsu Higher Education Institutions of China (Grant No. 13KJ430007), and the Qing Lan Project. HPCAT operations were supported by DOE-NNSA under Award No. DE-NA0001974 and DOE-BES under Award No. DE-FG02-99ER45775, with partial instrumentation funding by NSF. The gas loading was performed at GeoSoilEnviroCARS, APS, ANL, supported by Grants No. EAR-1128799 and No. DE-FG02-94ER14466. APS was supported by DOE-BES, under Contract No. DE-AC02-06CH11357. The neutron experiments were performed at the Swiss spallation neutron source SINQ, Paul Scherrer Institute, Villigen, Switzerland.


